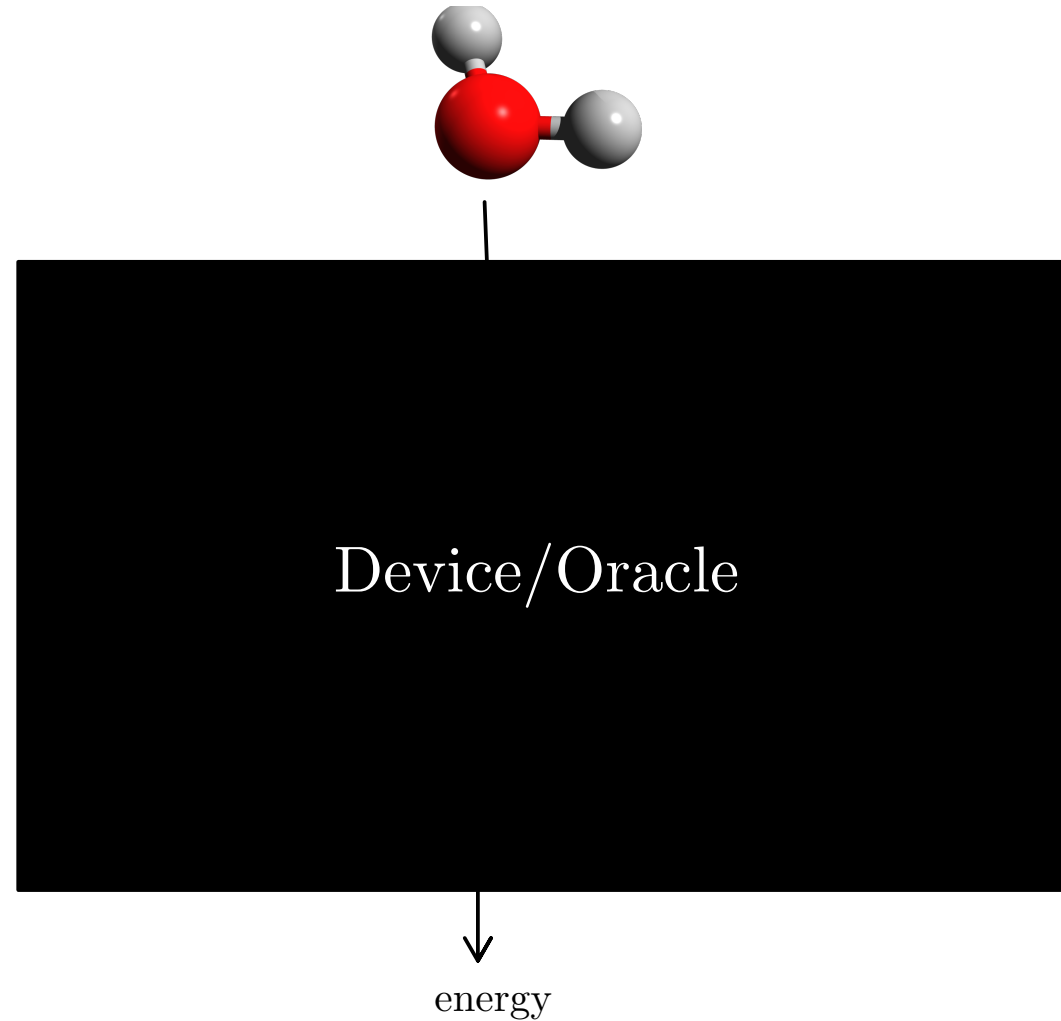


# Quantum Algorithms for Chemistry, Physics and Beyond

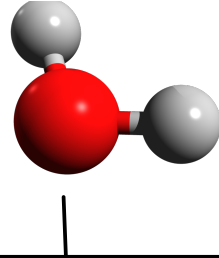
Jakob S. Kottmann



# Quantum Chemistry in a Nutshell



# Quantum Chemistry in a Nutshell

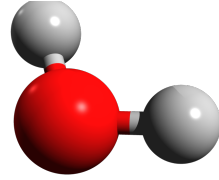


→ coordinates and charges →  $\{R_k, Z_k\}$  →  $V(r) = - \sum_k \frac{Z_k}{\|r - R_k\|}$

Device/Oracle

↓  
energy

# Quantum Chemistry in a Nutshell



→ coordinates and charges →  $\{R_k, Z_k\}$  →  $V(r) = - \sum_k \frac{Z_k}{\|r - R_k\|}$

real-space Hamiltonian

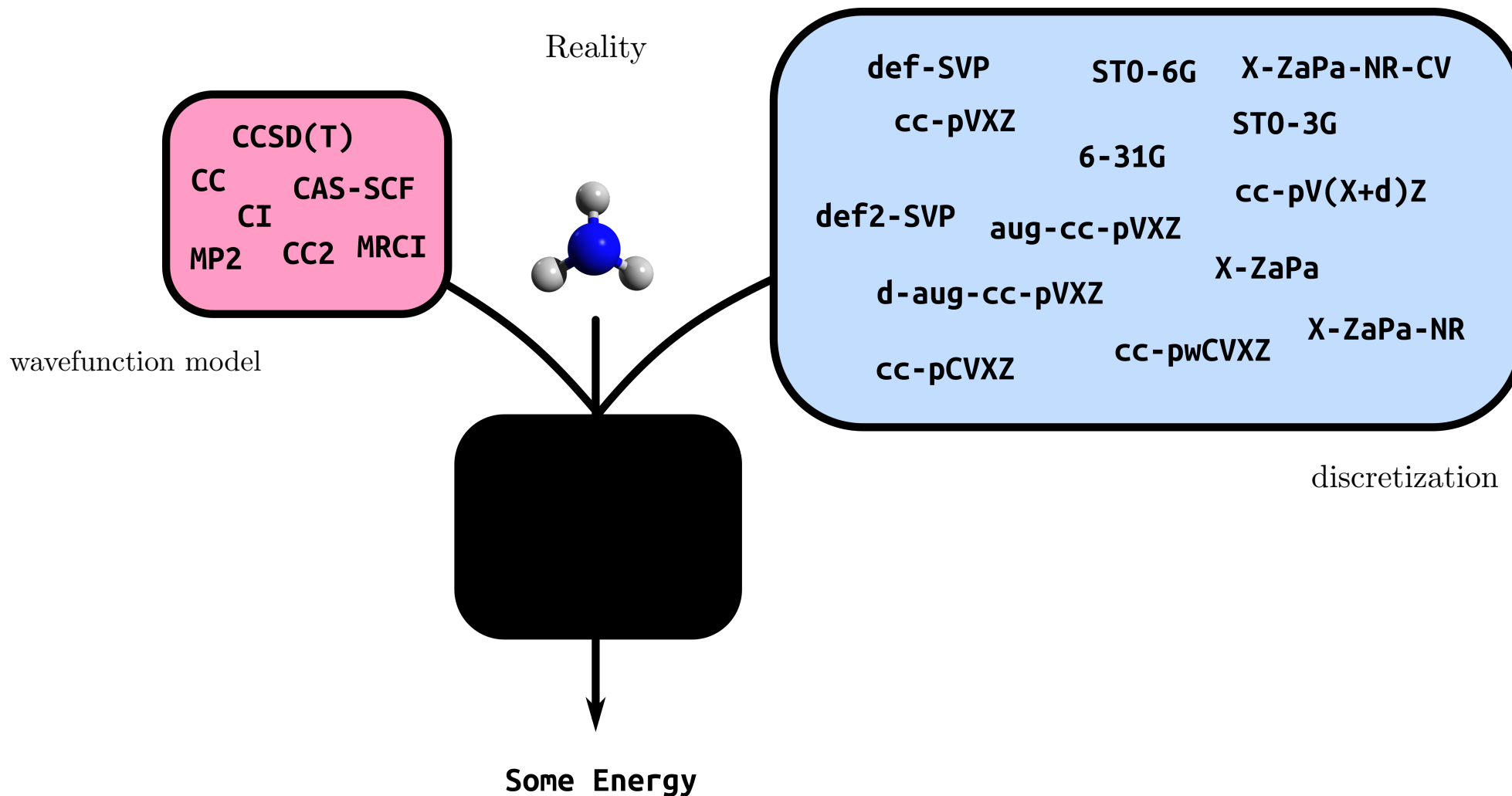
$$H(r_0, \dots, r_9) = - \sum_i \left( \frac{\nabla_{r_i}^2}{2} + V(r_i) \right) + \sum_{i < j} \frac{1}{\|r_i - r_j\|}$$

discretize

wavefunction model

energy

# Quantum Chemistry in a Nutshell



# Many-Body Physics: Splitting the Task

N-Body Problem

$$H(x, y) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \frac{1}{2} \frac{\partial^2}{\partial y^2} + \frac{1}{\|x - y\|} + f(x) + f(y)$$

Example: 2-Body Hamiltonian

# Many-Body Physics: Splitting the Task

N-Body Problem

$$H(x, y) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \frac{1}{2} \frac{\partial^2}{\partial y^2} + \frac{1}{\|x - y\|} + f(x) + f(y)$$

Example: 2-Body Hamiltonian

Classical Domain

effective one/two-body problems

$$F(x) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} + V(x, \phi)$$

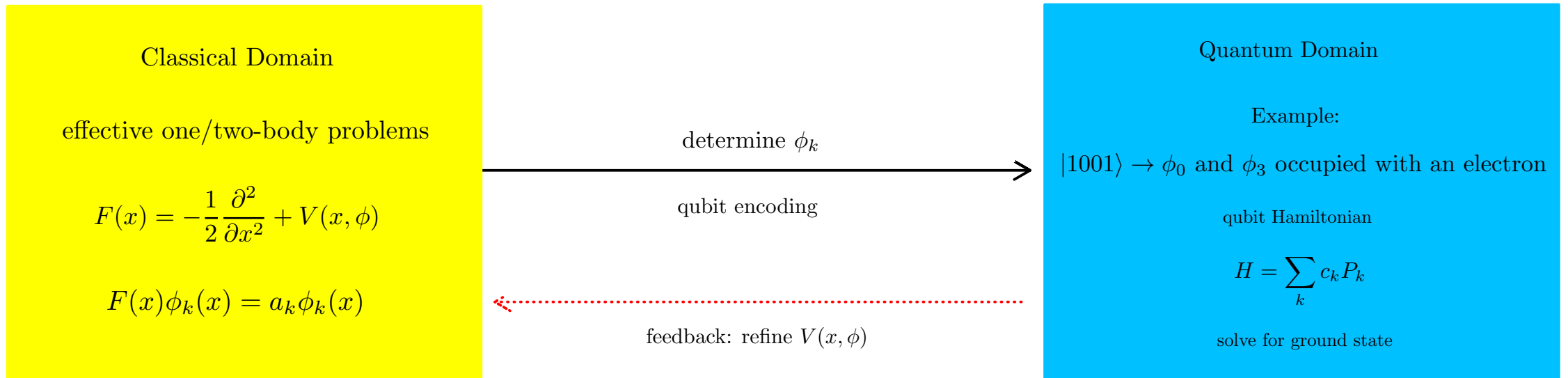
$$F(x)\phi_k(x) = a_k \phi_k(x)$$

# Many-Body Physics: Splitting the Task

N-Body Problem

$$H(x, y) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} - \frac{1}{2} \frac{\partial^2}{\partial y^2} + \frac{1}{\|x - y\|} + f(x) + f(y)$$

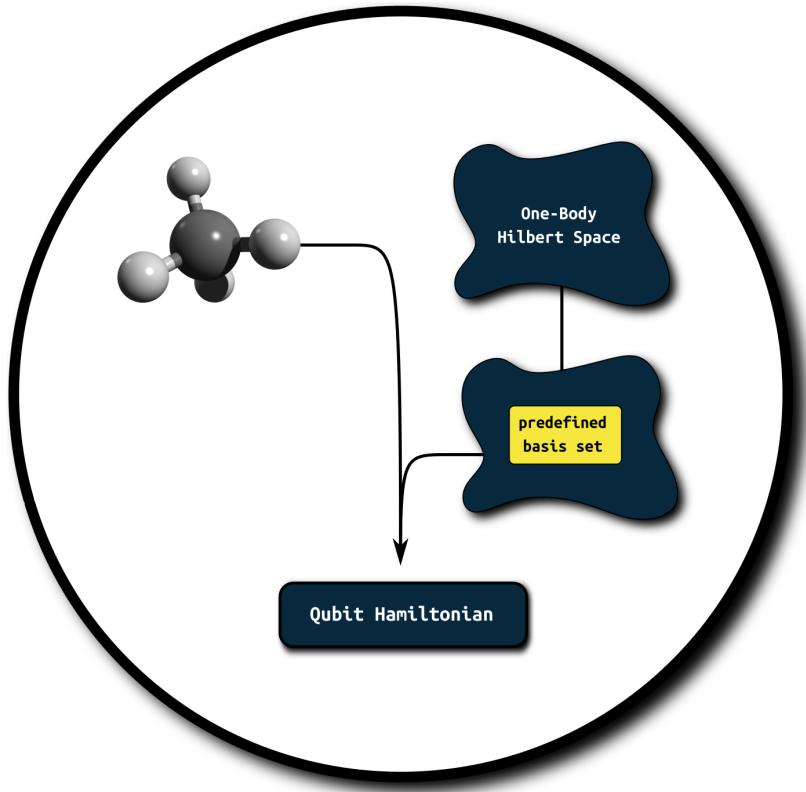
Example: 2-Body Hamiltonian





# Many-Body Physics: Splitting the Task

## Traditional Approach



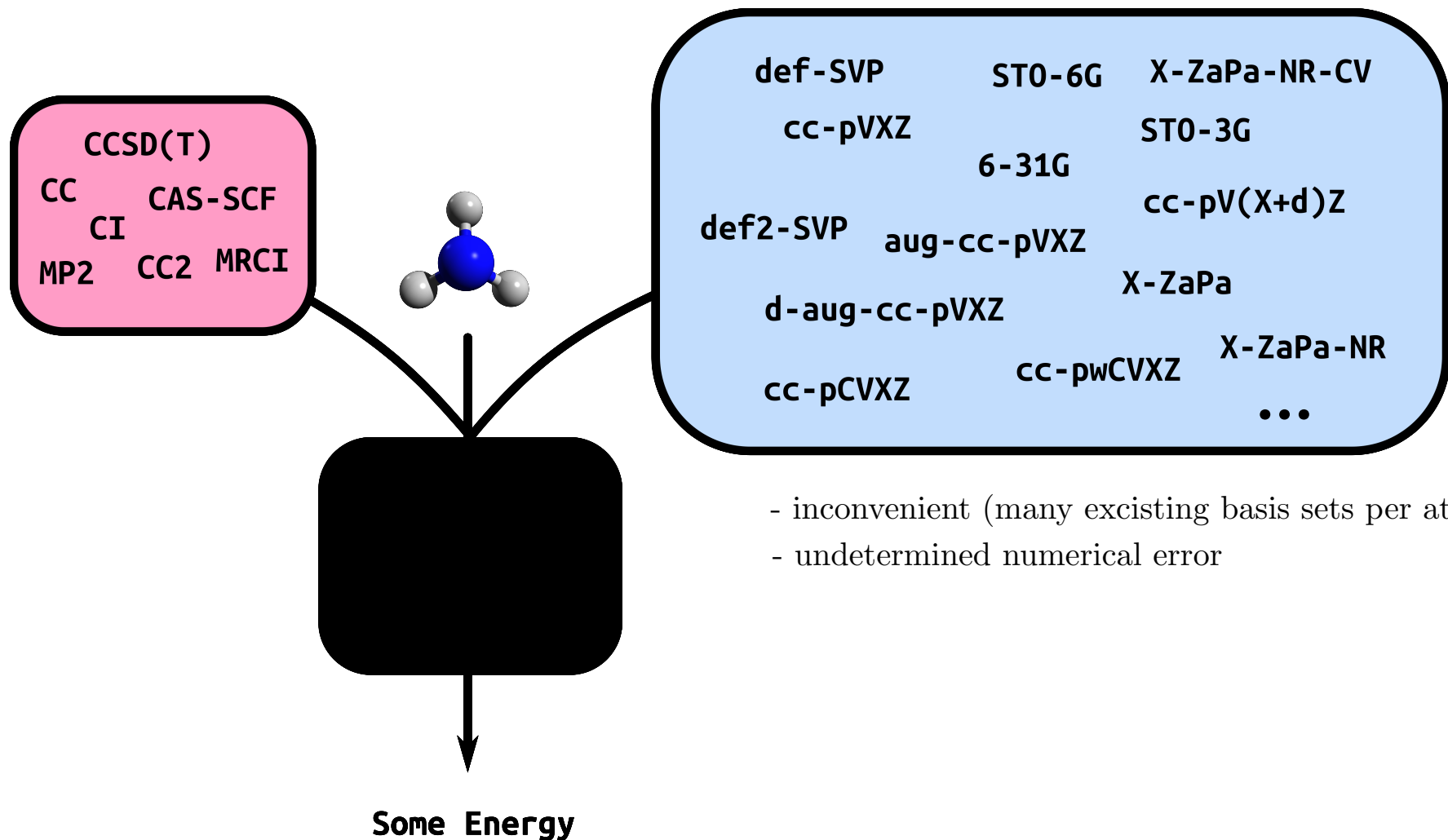
## Advantages

- well established
- fast integrals

## Drawbacks

- no black-box
- inconvenient (many existing basis sets per atom)
- more qubits necessary
- undetermined numerical error

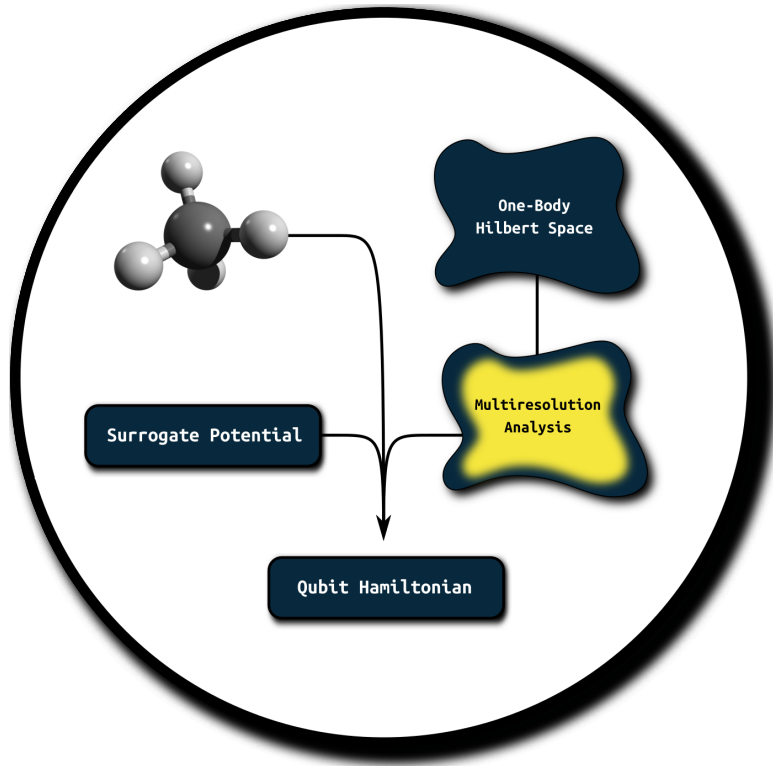
# Many-Body Physics: Splitting the Task



- inconvenient (many existing basis sets per atom)
- undetermined numerical error

# Many-Body Physics: Splitting the Task

## System-Adapted Approach



## Advantages

- defined numerical error
- can be treated as black-box
- low qubit numbers

## Drawbacks

- not well established
- comparably high classical computational cost  
formal scaling is often better though

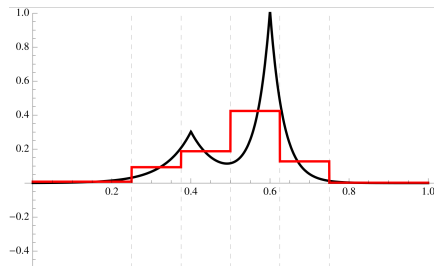
high-level blog article:

<https://aspuru.substack.com/p/bits-are-cheap-and-qubits-expensive>

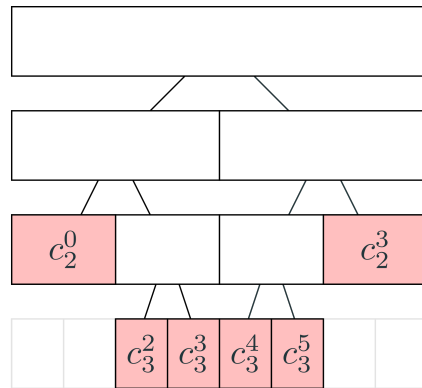
JSK, Schleich, Tamayo-Mendoza, Aspuru-Guzik. J.Chem.Phys.Lett. 2021

# System Adapted Approach: Behind the Scenes

## 1 Dimensional Example



$$|f\rangle = \sum_{nl} c_n^l |\varphi_n^l\rangle$$

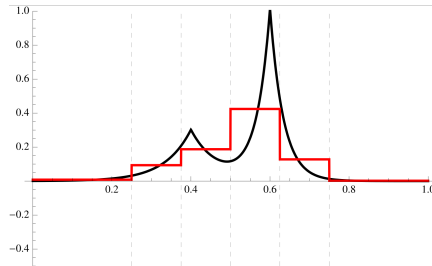


**github/madness**

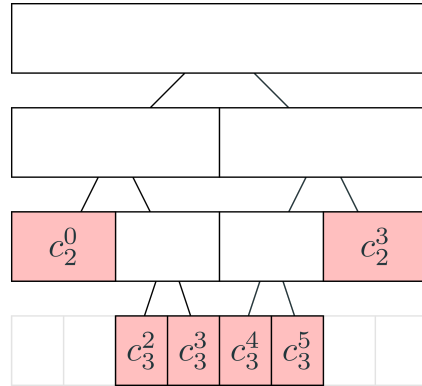
Harrison et. al.

# System Adapted Approach: Behind the Scenes

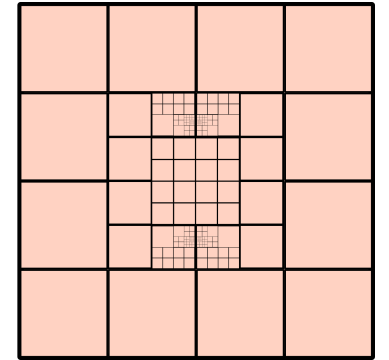
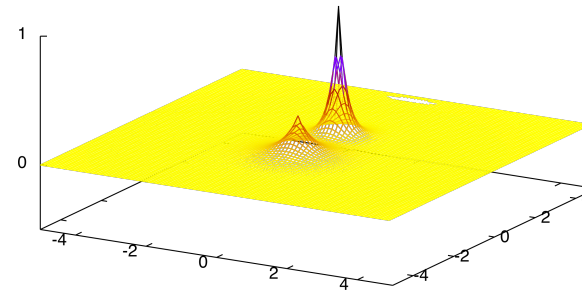
## 1 Dimensional Example



$$|f\rangle = \sum_{nl} c_n^l |\varphi_n^l\rangle$$



## 2 Dimensional Example

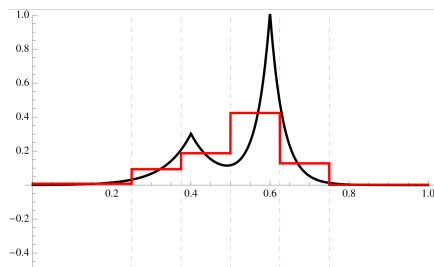


**github/madness**

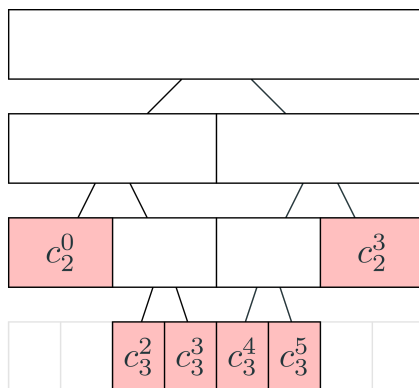
Harrison et. al.

# System Adapted Approach: Behind the Scenes

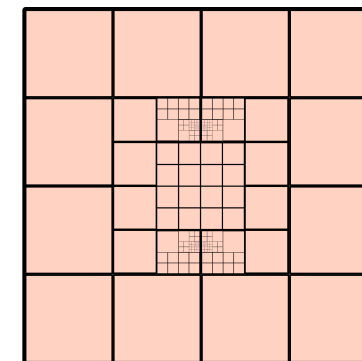
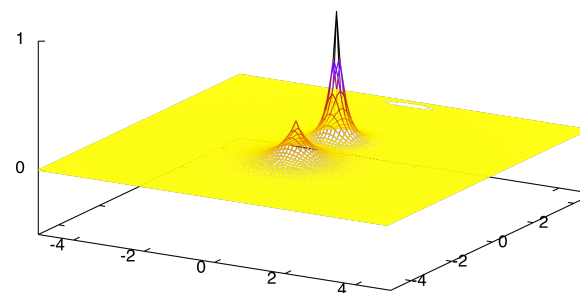
## 1 Dimensional Example



$$|f\rangle = \sum_{nl} c_n^l |\varphi_n^l\rangle$$



## 2 Dimensional Example



**github/madness**

Harrison et. al.

Classical Domain

effective one/two-body problems

$$F(x) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} + V(x, \phi)$$

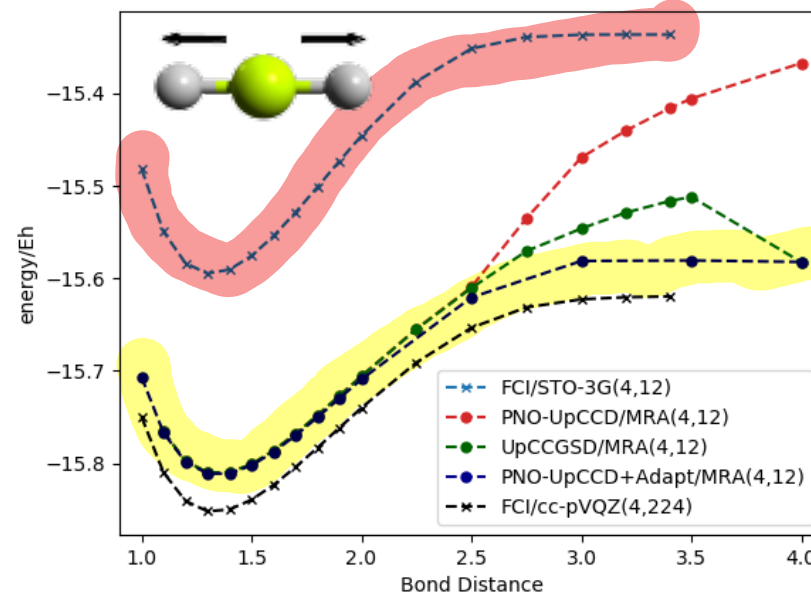
$$F(x)\phi_k(x) = a_k \phi_k(x)$$

# Basis-Set-Free Quantum Chemistry: Performance

System-Adapted Approach

System	Metric	MRA	GBS
He	MAX	4	4-10
Be	MAX	10	10-18
H <sub>2</sub>	NPE	4	20-56
H <sub>2</sub>	NPE	8	20-56
H <sub>2</sub>	NPE	20	56-120
H <sub>2</sub>	MAX	4	8
H <sub>2</sub>	MAX	8	20-56
H <sub>2</sub>	MAX	20	56
LiH	NPE	12-22	38-88
LiH	MAX	12	38-88
LiH	MAX	22	170-288
BH	NPE	12-22	38-88
BH	MAX	12-22	38-88
BeH <sub>2</sub>	NPE	12	46-114
BeH <sub>2</sub>	MAX	12	24-46
NH <sub>3</sub>	$\Delta E$	12-18	58-100

Traditional Approach



high-level blog article:

<https://aspuru.substack.com/p/bits-are-cheap-and-qubits-expensive>

JSK, Schleich, Tamayo-Mendoza, Aspuru-Guzik. J.Chem.Phys.Lett. 2021

what about the wavefunction model?

one way forward: quantum circuits



# Quantum Circuits

Task: Implement a unitary evolution

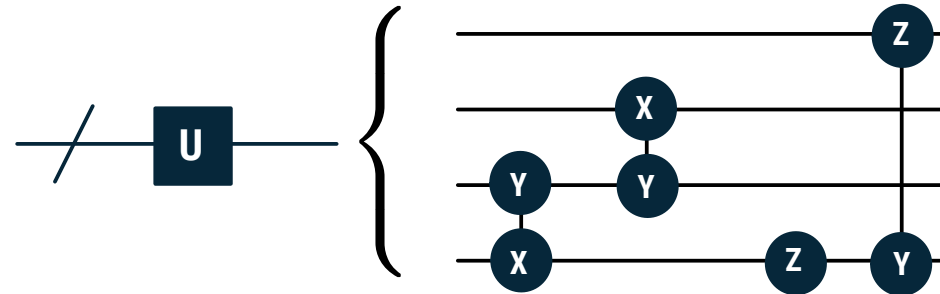
$$U = e^{i\theta G}$$

$G$  is Hermitian 1 or 2 qubit operator

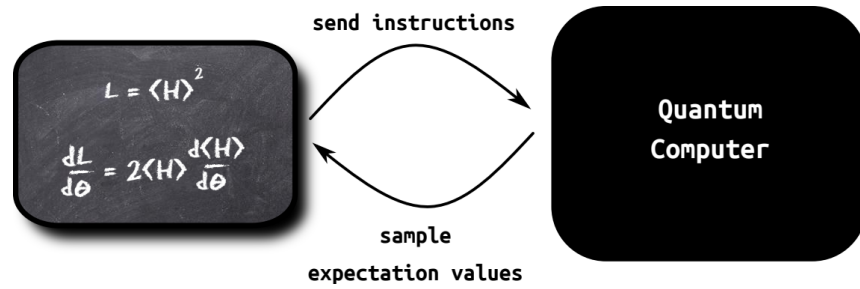
$\theta \in \mathbb{R}$  free parameter

Can be automatically differentiated

pioneers: M. Schuld *et.al* PRA 2019, PennyLane

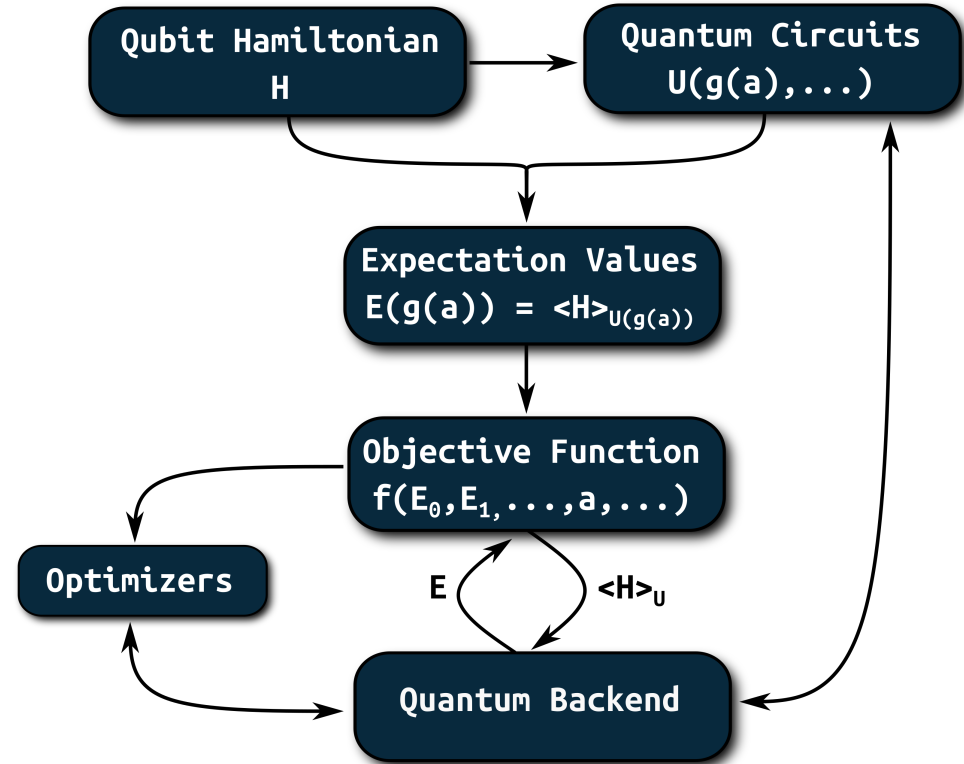


# Tequila: High Level Environment

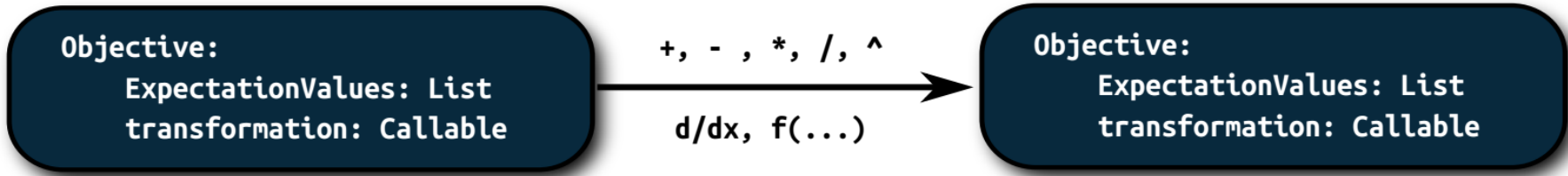


[github.com/aspuru-guzik-group/tequila](https://github.com/aspuru-guzik-group/tequila)

API inspired by madness library



concept



concept

**Objective:**  
**ExpectationValues: List**  
**transformation: Callable**

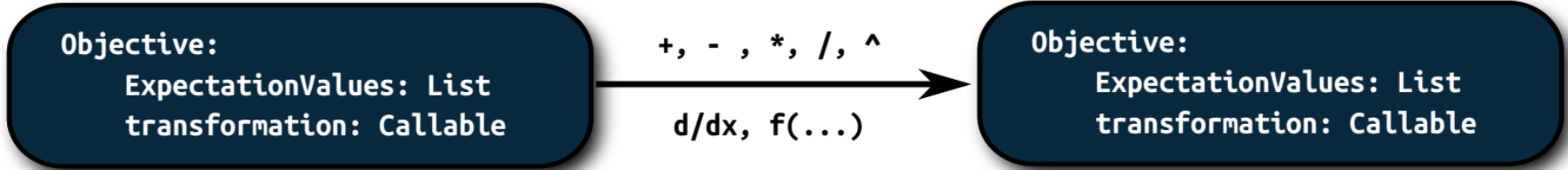
**+, - , \*, /, ^**  
**d/dx, f(...)**

**Objective:**  
**ExpectationValues: List**  
**transformation: Callable**

Example:  
high level code

```
01 = E0 + E1  
02 = 0.5*E0**2  
03 = 01**02
```

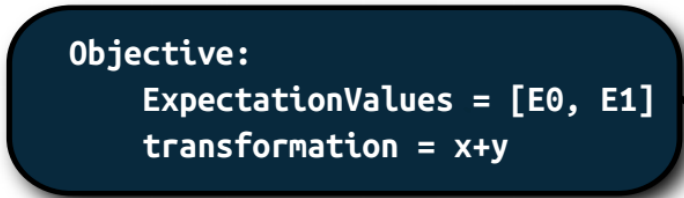
concept



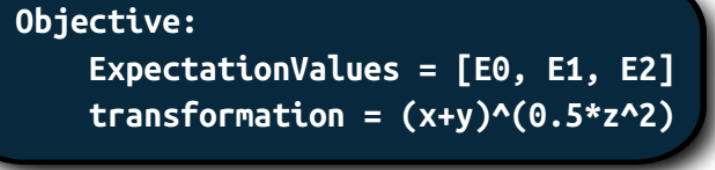
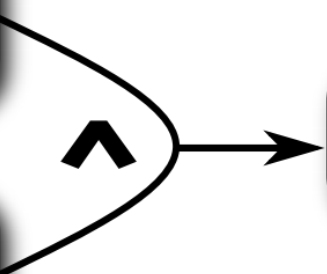
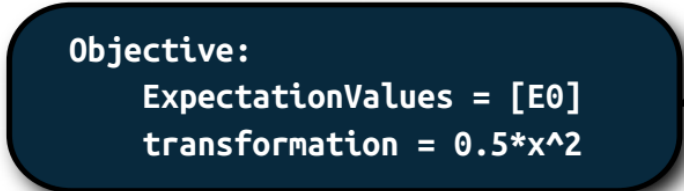
Example:  
high level code

```
O1 = E0 + E1  
O2 = 0.5*E0**2  
O3 = O1**O2
```

O<sub>1</sub>




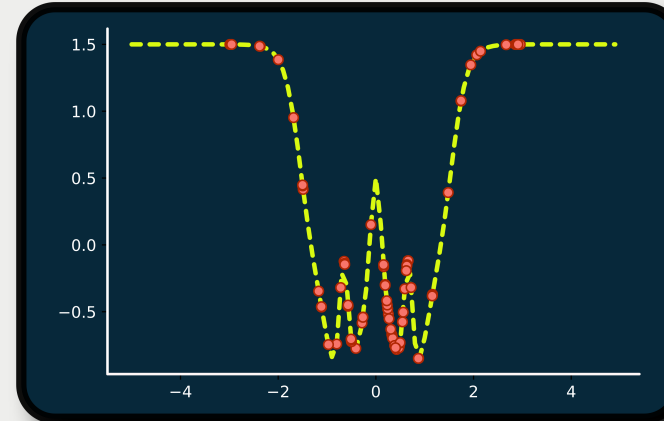
O<sub>2</sub>



O<sub>3</sub>

# Tequila: High Level Environment

$$H = -X(0)X(1) + \frac{1}{2}Z(0) + Y(1)$$

$$G = e^{-i\frac{t}{2}e^{-a^2}Y}$$
$$L = \langle H \rangle_{U(a)} + e^{-\left(\frac{d}{da}\langle H \rangle_{U(a)}\right)^2}$$



```
a = tq.Variable("a")
U = tq.gates.Ry(angle=(-a**2).apply(tq.numpy.exp)*pi, target=0)
U += tq.gates.X(target=1, control=0)
H = tq.QubitHamiltonian.from_string("-1.0*X(0)X(1)+0.5Z(0)+Y(1)")
E = tq.ExpectationValue(H=H, U=U)
dE = tq.grad(E, "a")
objective = E + (-dE**2).apply(tq.numpy.exp)
result = tq.minimize(method="phoenics", objective=objective)
```

## Some Development Examples

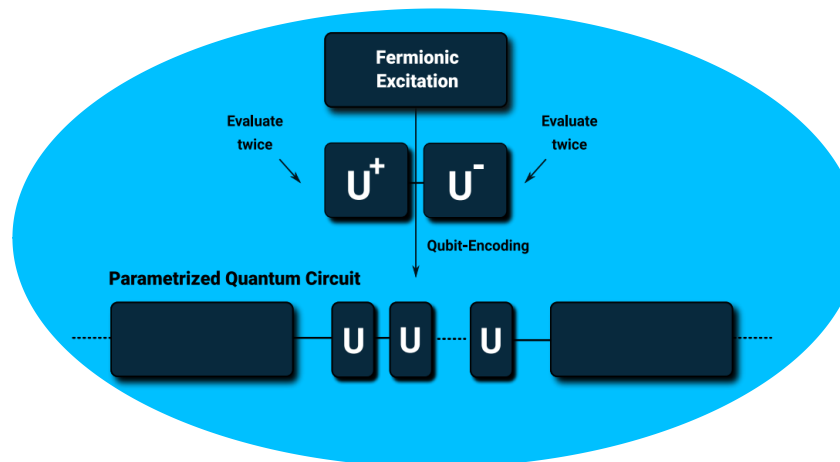
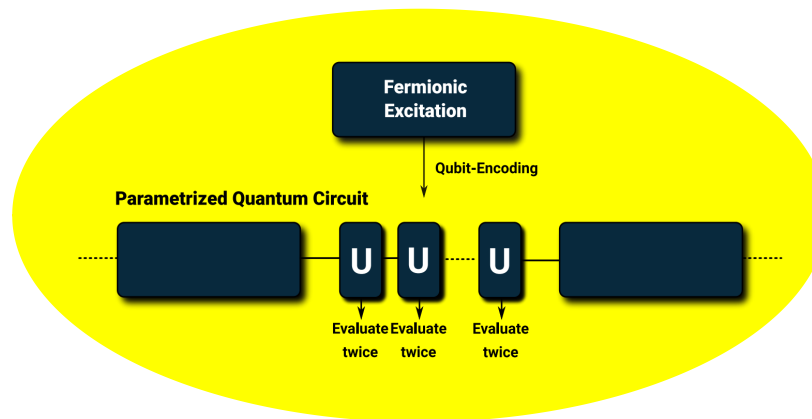
# Development Example: Advanced Gradients for Quantum Chemistry

gradient cost for  $n$  electron excitation

Generator Form	Gradient Cost	Strategy
$G_{\text{pq}} = \sum_i c_i \sigma_i$	$\mathcal{O}(2^{2n})$	shift-rule Eq. (6)
$G_{\text{pq}} = \frac{1}{2}(G_+ + G_-)$	4	fermionic-shift Eq. (16)
Real Wavefunctions		
$G_{\text{pq}} = \frac{1}{2}(G_+ + G_-)$	2	fermionic-shift Eq. (19)
Generator Approximation		
$G_{\text{pq}} \approx G_{\pm}$	2	shift-rule Eq. (6)

Basic building blocks for Unitary Coupled-Cluster

recent review: A. Anand *et.al.* 2021



→ generalizable

follow-ups:

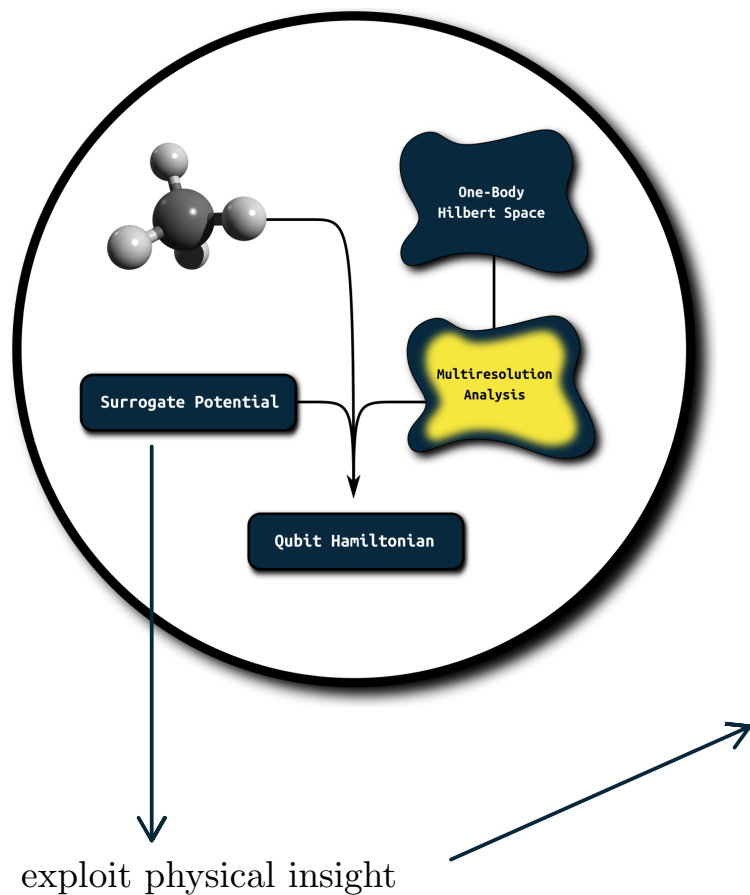
Izmaylov *et.al.* 2021

Anselmetti *et.al.* 2021

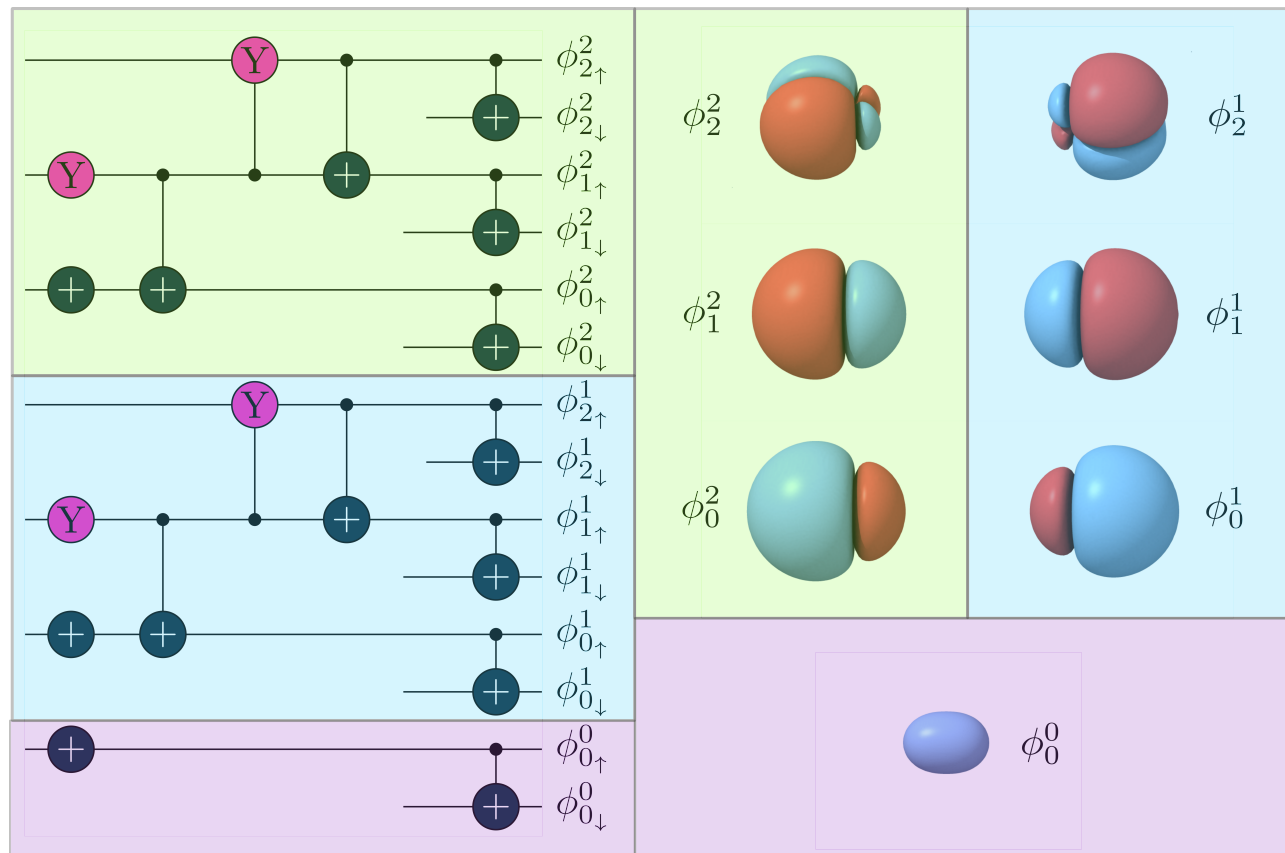
Wierichs *et.al.* 2021



# Development Example: Separable Pair Approximations



High level circuit design through physical principles



# Development Example: Separable Pair Approximations

classically simulatable

→ cheap initial states

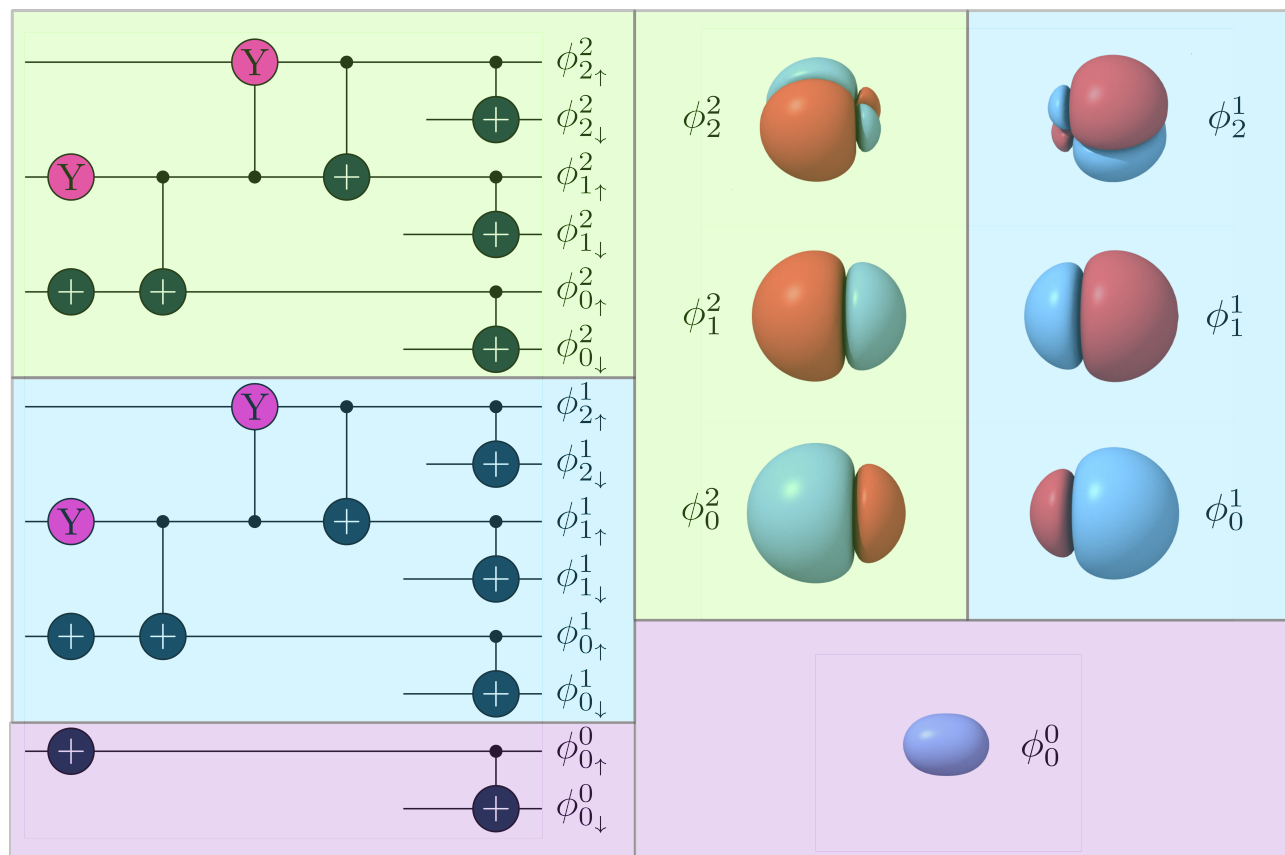
separated pairs

→ distributed schemes

→ hybrid simulation

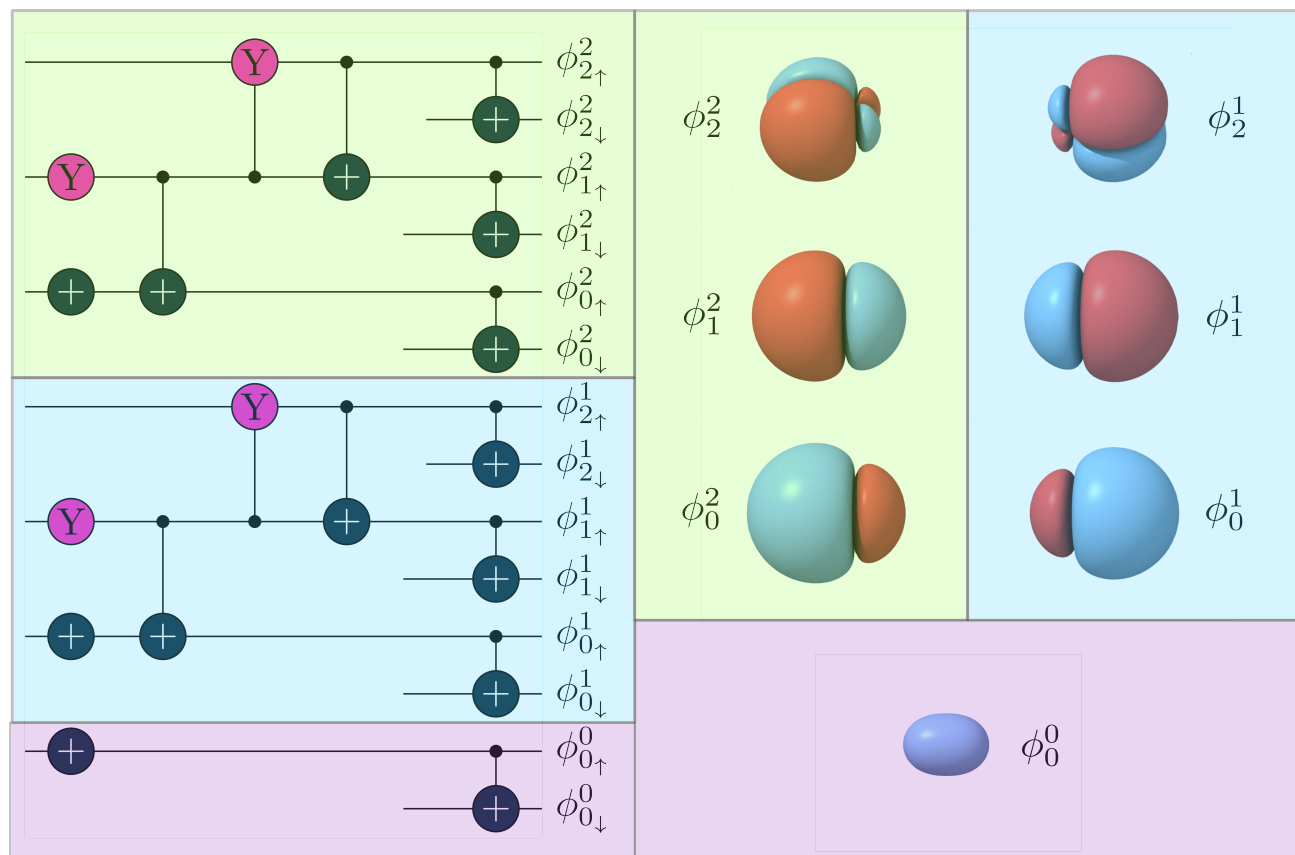
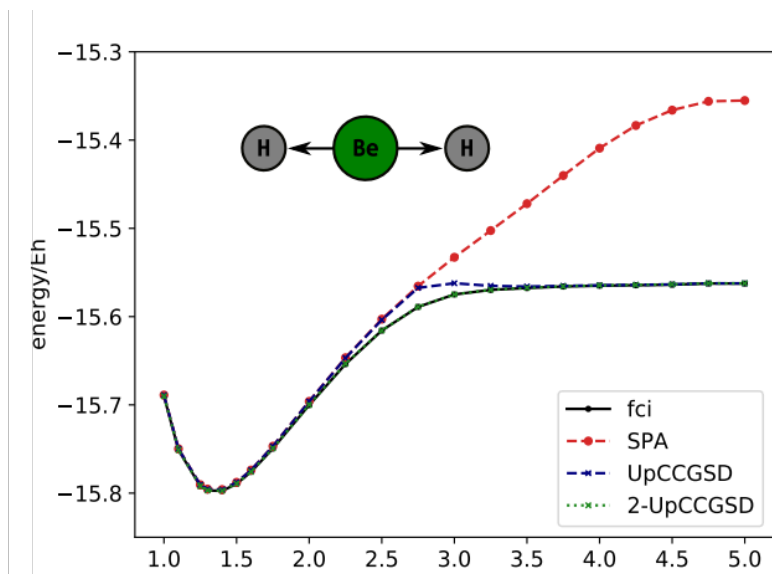
Molecule( $N_e, N_q$ )	$N_{\text{param}}$	$N_{\text{cnot}}$	Depth
H <sub>2</sub> (2,4)	1	3	3
LiH(2,10)	4	15	18
BeH <sub>2</sub> (4,8)	2	6	3
BeH <sub>2</sub> (6,14)	4	15	7
BH <sub>3</sub> (6,12)	3	9	3
N <sub>2</sub> (6,12)	3	9	3
C <sub>2</sub> H <sub>4</sub> (12,24)	6	18	3
H <sub>2</sub> O <sub>2</sub> (14,28)	7	21	3
C <sub>2</sub> H <sub>6</sub> (14,28)	7	21	3
C <sub>2</sub> H <sub>6</sub> (2,12)	5	19	23
C <sub>2</sub> H <sub>6</sub> (14,84)	35	133	23

High level circuit design through physical principles

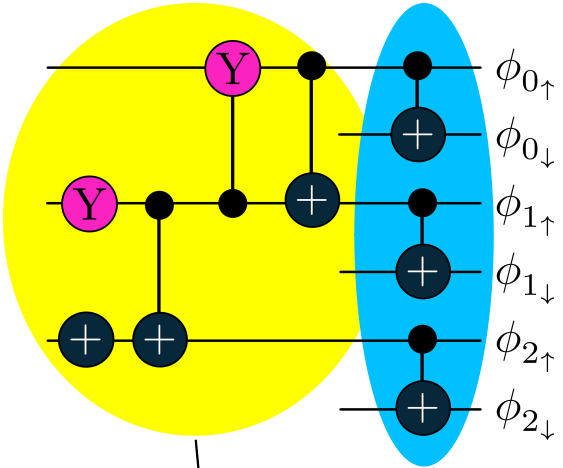


# Development Example: Separable Pair Approximations

High level circuit design through physical principles



# Possible Extensions



combine with other approaches:  
 Anselmetti, arXiv:2104.05695, 2021

```
class JordanWigner(EncodingBase):
    ...
    def hcb_to_me(self, *args, **kwargs):
        U = QCircuit()
        for i in range(self.n_orbitals):
            U += X(target=self.down(i), control=self.up(i))
        return U
```

$$U_{HCB}^X = U_{JW}^X U_{HCB}^{JW}$$

$X \in \{\text{Bravyi-Kitaev}, \dots\}$

quasi-local codes: Chien & Whitfield, arXiv:2009.11860  
 BKSF: Setia, Bravyi, Mezzacapo, Whitfield, PRR 2019

transfer to Jordan-Wigner

wavefunction preparation in paired sector

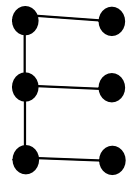
"Pairing models", "Hardcore-Boson", "Seniority-Zero"

access to similar ideas:

Elfving *et.al*, PRA, 2021

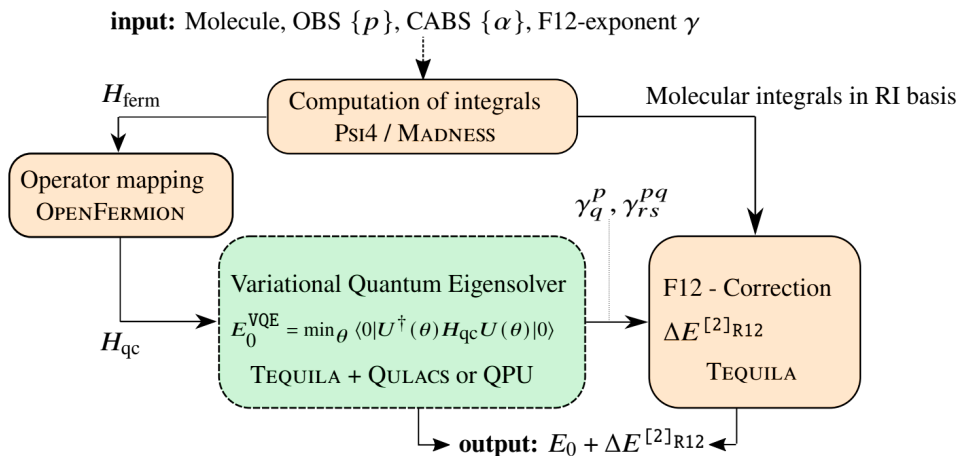
Khamoshi, Evangelista, Scuseria, QST, 2020

explore qubit connectivity



# Recent Developments

Explicit Correlation



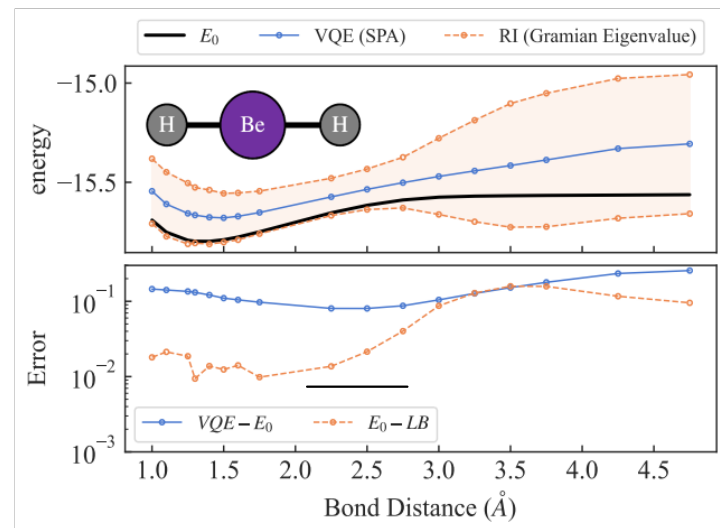
P. Schleich, JSK, A. Aspuru-Guzik, Arxiv:2110.06812, 2021

See also: Master thesis from Philipp Schleich (detailed introduction)

Robustness Intervals: Weber *et.al*, arXiv:2110.09793

	SDP		Gramian	
	Expectation $\langle A \rangle_\sigma$		Expectation $\langle A \rangle_\sigma$	Eigenvalue $\lambda$
Lower Bound	$(1 - 2\epsilon)\langle A \rangle_\rho - 2\sqrt{\epsilon(1 - \epsilon)(1 - \langle A \rangle_\rho^2)}$		$(1 - 2\epsilon)\langle A \rangle_\rho - 2\sqrt{\epsilon(1 - \epsilon)}\Delta A_\rho + \frac{\epsilon\langle A^2 \rangle_\rho}{\langle A \rangle_\rho}$	$\langle A \rangle_\rho - \Delta A_\rho\sqrt{\frac{\epsilon}{1 - \epsilon}}$
Upper Bound	$(1 - 2\epsilon)\langle A \rangle_\rho + 2\sqrt{\epsilon(1 - \epsilon)(1 - \langle A \rangle_\rho^2)}$		—	$\langle A \rangle_\rho + \Delta A_\rho\sqrt{\frac{\epsilon}{1 - \epsilon}}$
Assumptions	$-1 \leq A \leq 1$		$A \geq 0$	$\sigma =  \psi\rangle\langle\psi  \wedge A \psi\rangle = \lambda \psi\rangle$

TABLE I. Overview of bounds for the expectation values and eigenvalues of an Hermitian operator  $A$  under a target state  $\sigma$ , with  $\rho$  an approximation of  $\sigma$ . For the eigenvalue bound,  $\sigma = |\psi\rangle\langle\psi|$  is the density operator corresponding to the eigenstate  $|\psi\rangle$  with eigenvalue  $\lambda = \langle\psi|A|\psi\rangle$ . We remark that the SDP lower and upper bounds are valid for fidelities with  $\mathcal{F}(\rho, \sigma) \geq 1 - \epsilon$  for  $\epsilon \geq 0$  such that  $\epsilon \leq \frac{1}{2}(1 + \langle A \rangle_\rho)$  and  $\epsilon \leq \frac{1}{2}(1 - \langle A \rangle_\rho)$ , respectively. The Gramian lower bound for expectation values is valid for  $\epsilon \geq 0$  with  $\sqrt{1 - \epsilon}/\epsilon \geq \Delta A_\rho/\langle A \rangle_\rho$ .



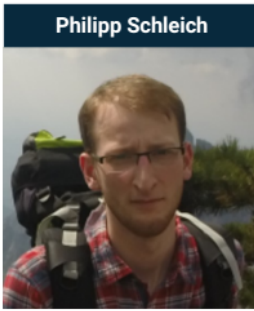
gradients, orbitals, circuits: All used in black-box fashion



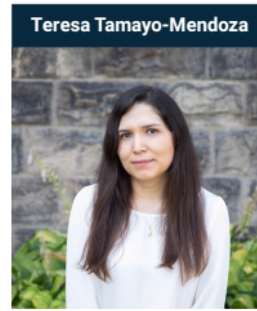
Alán Aspuru-Guzik



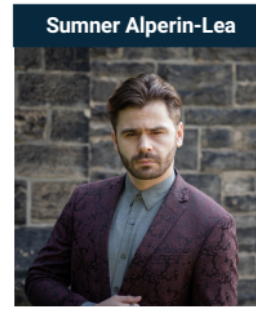
Abhinav Anand



Philipp Schleich



Teresa Tamayo-Mendoza



Sumner Alperin-Lea



Alba Cervera-Lierta

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Georgios Tsilimigkounakis,  
Claudia Zendejas-Morales,  
Tanya Garg

Maurice Weber, Arianne van den Griend

